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STD-R-1071
THE ELECTRICAL PROPERTIES
OF SEAWATER (INCLUDING
CONDUCTIVITY RELAXATION)

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STD-R-1071
THE ELECTRICAL PROPERTIES
OF SEAWATER (INCLUDING
CONDUCTIVITY RELAXATION)

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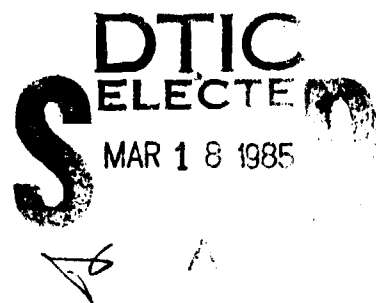
M. E. THOMAS

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**FINAL REPORT
TASK - ZEDO
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**SUBMARINE TECHNOLOGY DEPARTMENT
The Johns Hopkins University
Applied Physics Laboratory**



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SECURITY CLASSIFICATION OF THIS PAGE

REPORT DOCUMENTATION PAGE

1a. REPORT SECURITY CLASSIFICATION Unclassified			1b. RESTRICTIVE MARKINGS	
2a. SECURITY CLASSIFICATION AUTHORITY			3. DISTRIBUTION / AVAILABILITY OF REPORT	
2b. DECLASSIFICATION / DOWNGRADING SCHEDULE				
4. PERFORMING ORGANIZATION REPORT NUMBER(S) STD-R-1071			5. MONITORING ORGANIZATION REPORT NUMBER(S)	
6a. NAME OF PERFORMING ORGANIZATION Applied Physics Lab Johns Hopkins University		6b. OFFICE SYMBOL (If applicable)	7a. NAME OF MONITORING ORGANIZATION Office of Naval Research, Code 422CS	
6c. ADDRESS (City, State, and ZIP Code) Johns Hopkins Road Laurel, MD 20707			7b. ADDRESS (City, State, and ZIP Code) Arlington, VA 22218	
8a. NAME OF FUNDING / SPONSORING ORGANIZATION ONR		8b. OFFICE SYMBOL (If applicable) Code 420	9. PROCUREMENT INSTRUMENT IDENTIFICATION NUMBER	
8c. ADDRESS (City, State, and ZIP Code) Arlington, VA 22218			10. SOURCE OF FUNDING NUMBERS	
			PROGRAM ELEMENT NO.	PROJECT NO.
			TASK NO.	WORK UNIT ACCESSION NO.
11. TITLE (Include Security Classification) The Electrical Properties of Seawater (Including Conductivity Relaxation)				
12. PERSONAL AUTHOR(S) M. E. Thomas				
13a. TYPE OF REPORT Technical Report		13b. TIME COVERED FROM 10/83 TO 9/84		14. DATE OF REPORT (Year, Month, Day) 1984, July
15. PAGE COUNT 34				
16. SUPPLEMENTARY NOTATION Approved for Public Release, Distribution Unlimited				
17. COSATI CODES			18. SUBJECT TERMS (Continue on reverse if necessary and identify by block number)	
FIELD	GROUP	SUB-GROUP	(a) Dielectrical Properties (c) Conductivity	
			above one GHz, and Relaxation	
			(b) Seawater (d) Microwave. A	
19. ABSTRACT (Continue on reverse if necessary and identify by block number) The electrical properties of seawater for ac currents or fields of 1 GHz or higher frequencies are examined with particular emphasis on the effects of relaxation of the charge carriers with the frequency of the applied field. Such effects have not been definitively quantified either experimentally or theoretically. New experiments are conducted but are not conclusive. It is possible that the conductivity relaxation time is similar to the Debye relaxation time making direct observations of this effect difficult. <i>Keynote included</i>				
20. DISTRIBUTION / AVAILABILITY OF ABSTRACT <input checked="" type="checkbox"/> UNCLASSIFIED/UNLIMITED <input type="checkbox"/> SAME AS RPT. <input type="checkbox"/> DTIC USERS			21. ABSTRACT SECURITY CLASSIFICATION Unclassified	
22a. NAME OF RESPONSIBLE INDIVIDUAL W. M. Chambers			22b. TELEPHONE (Include Area Code) 301/953-5123	22c. OFFICE SYMBOL n/a

DD FORM 1473, 84 MAR

83 APR edition may be used until exhausted.
All other editions are obsolete.

SECURITY CLASSIFICATION OF THIS PAGE

UNCLASSIFIED

The Johns Hopkins University
APPLIED PHYSICS LABORATORY
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1. INTRODUCTION

A better understanding of the electromagnetic properties of aqueous electrolytes has important implications for oceanography and biology. The purpose of this study is to consider the conductivity relaxation time of seawater. Potentially, this phenomenon influences the electrical properties of seawater above 5GHz. Also, the dielectric resonant bands in the far-infrared affect the electrical properties beyond 5 GHz, and their effects are considered based on available information in the literature.

Although the study of aqueous electrolytic solutions has a long history, much yet remains to be understood. Essentially, the electrical properties of solutions of high concentration [>0.1 Molarity (M)] (i.e., seawater) are not completely understood and are typically represented by empirical models. To improve the present understanding of seawater conductivity, the determination of its relaxation time is important because it gives information about the mobility of the charge carriers.

To see how the properties of mobility influence the impedance of electrolytes, consider the standard definition for conductivity (Reference 1)

$$\sigma = \sum_{i=1}^N (|e_{+i}| n_{+i} \mu_{+i} + |e_{-i}| n_{-i} \mu_{-i}); \sum_{i=1}^N (e_{+i} n_{+i} + e_{-i} n_{-i}) = 0 \quad (1)$$

where $n_{+,-i}$ is the number density of the i th $+,-$ charge carrier, $e_{+,-i}$ is the charge of the i th $+,-$ charge carrier, $\mu_{+,-i}$ is the carrier mobility of the i th $+,-$ charge carrier, and N is the number of solutes. Since most electrolytes composing the ocean are 1:1, it is assumed that

$$|e_{+i}| n_{+i} = |e_{-i}| n_{-i} = \rho_i$$

The conductivity for a single 1:1 solute simplifies to

$$\sigma = \rho (\mu_+ + \mu_-) . \quad (2)$$

A simple expression for the carrier mobility can be obtained from the following equations (Reference 2)

$$\mu = \frac{\langle v \rangle}{E}$$

and

$$m \frac{d\langle v \rangle}{dt} + mg \langle v \rangle = eE . \quad (3)$$

where $\langle v \rangle$ is the mean velocity of the charge carrier, m is the mass, g is the damping constant, e is the charge of the charge carrier, and E is the applied electric field. Let $\langle v \rangle$ and E be time harmonic as $e^{j\omega t}$. The solution of Equation (3) results in the following

$$\mu = \frac{\mu_0}{1 + j \frac{\omega}{g}} ; \mu_0 = \frac{e}{mg} . \quad (4)$$

Substituting this result into Equation (2), the conductivity becomes

$$\sigma = \rho \left[\frac{\mu_{0+}}{1 + j \frac{\omega}{g_+}} + \frac{\mu_{0-}}{1 + j \frac{\omega}{g_-}} \right] . \quad (5)$$

It is assumed that the damping forces will be similar for both the positive and negative charge carriers. Therefore,

$$g_+ = g_- = g .$$

The "g" parameter is the reciprocal of the conductivity relaxation time. This interpretation is easily seen from Equation (3). When the applied E-field is turned off, the charge carriers relax back to equilibrium and the mean velocity becomes

$$\langle v \rangle = \langle v(t=0) \rangle e^{-tg}.$$

The conductivity relaxation time, $1/g$, is designated " τ_c ." Thus,

$$\begin{aligned}\sigma &= \rho(\mu_{o+} + \mu_{o-}) \frac{1}{1 + j \frac{\omega}{g}} \\ &= \sigma_o \frac{1}{1 + j \omega \tau_c}.\end{aligned}$$

where τ_c is the conductivity relaxation time related to mobility. The factor σ_o represents the conductivity according to the Debye-Falkenhagen theory (Reference 3) and the remaining factor represents the correction for inertia.

The impedance can now be determined for a parallel plate probe as

$$\begin{aligned}Z &= \frac{l}{A\sigma} - j \frac{1}{\omega C} = \frac{l}{A\sigma} \left(1 + j \frac{\omega}{g}\right) - j \frac{1}{\omega C} \\ &= R_o + j \left(\omega L - \frac{1}{\omega C}\right); \quad L = \frac{l}{A\sigma g} = \frac{R_o}{g}\end{aligned}\tag{7}$$

where l is the length of the sample and A is the cross-sectional area of the plate. R_0 is the DC resistance, L is the effective inductance, and C is the capacitance of the solution between the plates. Since the ions cannot move instantaneously with the applied field, the current will lag the voltage. The conductivity relaxation time, τ_0 , is determined by computing L/R_0 .

Impedance measurements of this type should be used instead of microwave cavity measurements, which are influenced by both the conductive and dielectric properties of the measured medium. By keeping the frequency of the applied field low enough (< 1 GHz) so that the real part of the dielectric constant is constant, the imaginary part is very small and the conductivity high enough so that the capacitive (i.e., dielectric) character is shorted out; then measurement can readily be made. The main problem remaining is minimizing the stray capacitance and residual inductance.

Section 2 describes the impedance experiment designed to measure the conductivity relaxation time of seawater and related electrolytes. Section 3 reviews the present model used for describing the electrical properties of seawater and discusses the effect of conductivity relaxation and other mechanisms on this model.

2. ELECTRICAL IMPEDANCE MEASUREMENTS

The major thrust of this project is to measure the electrical impedance of seawater-related electrolytes. A special conductivity cell was designed and constructed for this purpose. The following describes the experimental apparatus used and presents the results of the experiment.

2.1

Experimental Apparatus

Figure 1 illustrates the experimental apparatus. The impedance measurements were made with a Hewlett-Packard (HP) multifrequency LCR meter model number 4275A, which operates at ten different frequencies between 10 kHz and 10 MHz. The LCR meter was computer controlled by an HP 9845A computer. A conductivity test cell was connected to the LCR meter in such a way that the shields of the 4-terminal network were connected to the shield enclosing the test cell. As shown in Figure 1, the current passing through the cell also was carried by the shield thus canceling the magnetic field generated by the central conductor. This reduced the residual inductance of the network.

Figure 2 shows the unique design of the modified Jones cell that was used as the test cell. The cell can be filled and drained without disconnecting it from the LCR meter. Also, the cell can be temperature controlled by an external bath. The cell constant is $17.2 (\pm 0.15) \text{cm}^{-1}$. The temperature of the cell is monitored by placing a probe inside a well in contact with the circulating bath.

The accuracy of the impedance measurements depends on the LCR meter and background effects of the test cell. The LCR meter does open- and short-circuit measurements of the test network and measures the stray capacitance and residual inductance. The values are stored and then subtracted from the actual measurements. The open-circuit measurement is done with the cell empty. The short-circuit measurement is done with the cell filled with mercury. In this way, most of the stray capacitance and residual inductance are accounted for. Without this capability these measurements would not be possible. The LCR meter reduces the stray capacitance and the residual inductance to some nominal level. These levels are measured before every experiment. Figure 3 shows the accuracy of the magnitude and phase of the impedance.

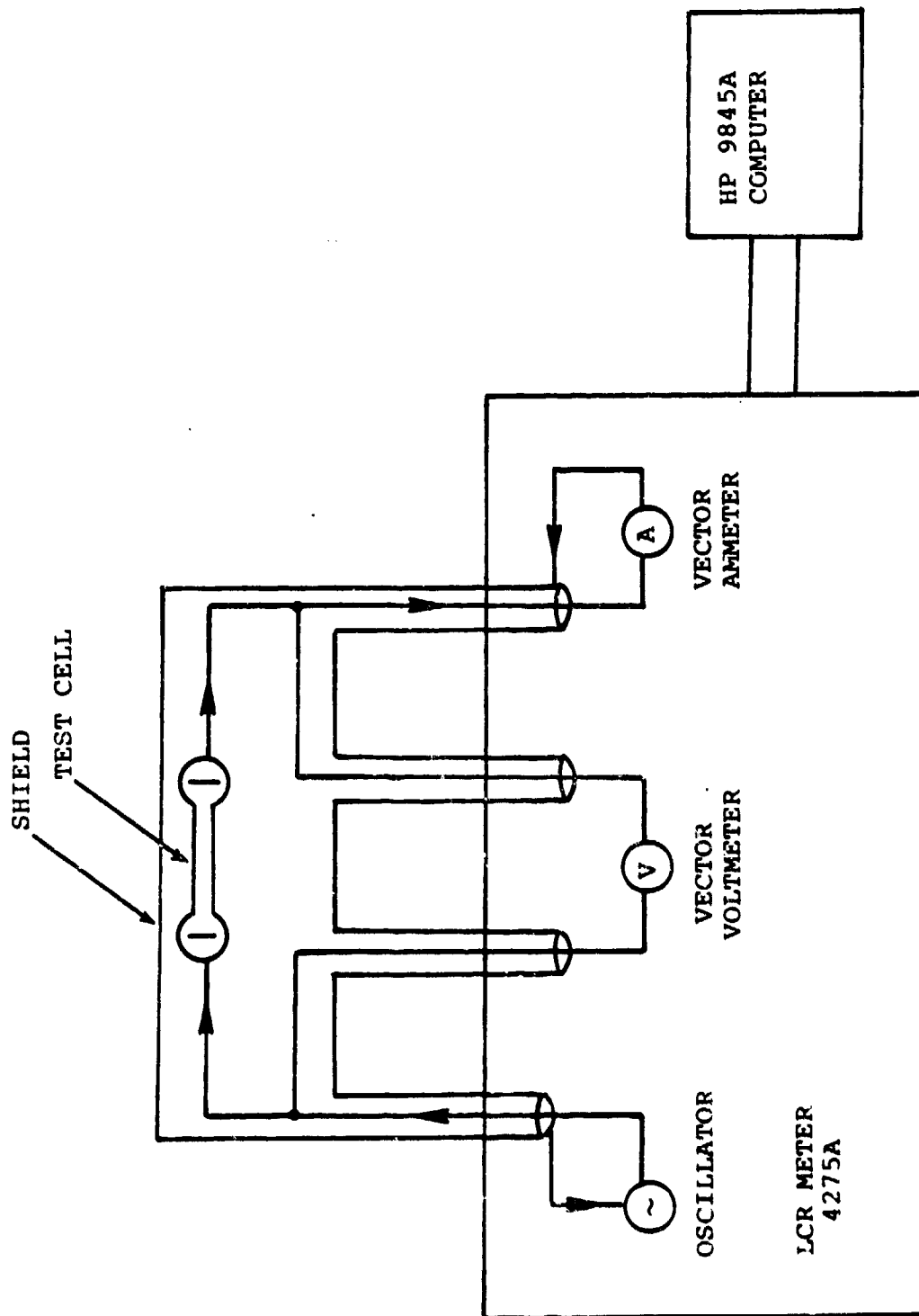


Figure 1 Experimental Apparatus (vector meters measure magnitude and phase)

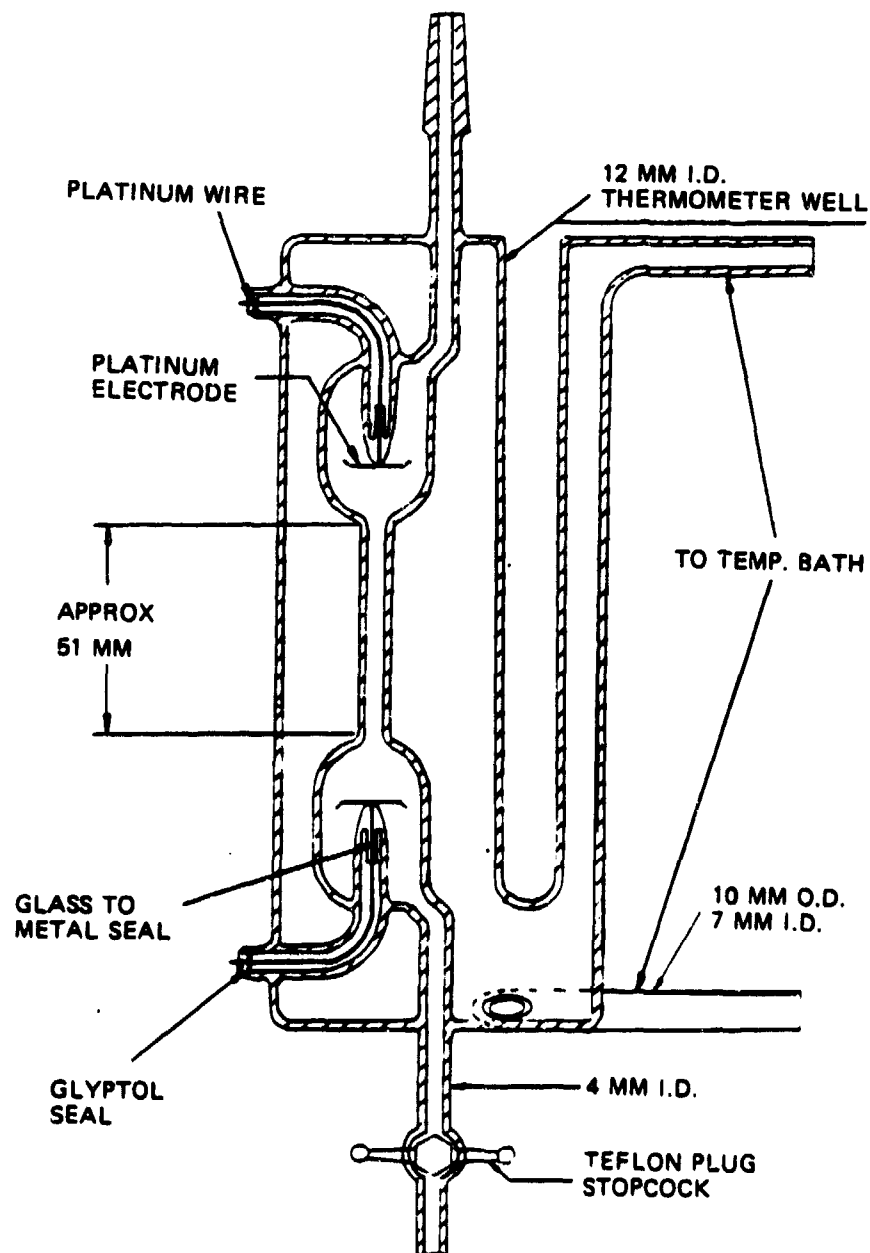


Figure 2 Modified Jones Test Cell

|Z| - θ MEASUREMENT

Z RANGE	TEST FREQUENCY																				
	10M Ω	1000k Ω	100k Ω	10k Ω	1000 Ω	100 Ω	10 Ω	1000m Ω	10kHz	20kHz	40kHz	100kHz	200kHz	400kHz	1MHz	2MHz	4MHz	10MHz			
	$\frac{25 \pm 1}{0.1'' \pm 0.1''/e}$		$\frac{25 \pm 1}{0.5'' \pm 0.5''/e}$																		

Equations in table represent:

Impedance accuracy
Phase angle accuracy

α , $1/\alpha$: See Figure A Accuracy Coefficient Graph.

θ measurement range:

-180.000° - +180.000°

Display counts for |Z| and θ (normal mode):

Ranges	Z	θ
	*36 - 1999	0 - 18000
	0 - 19999	0 - 18000

*Approximate value (unspecified).

Number of significant digits displayed for |Z| and θ depend on test signal level, range and frequency (5 digits max.).

Accuracies in lined areas are unspecified.

ACCURACY COEFFICIENTS

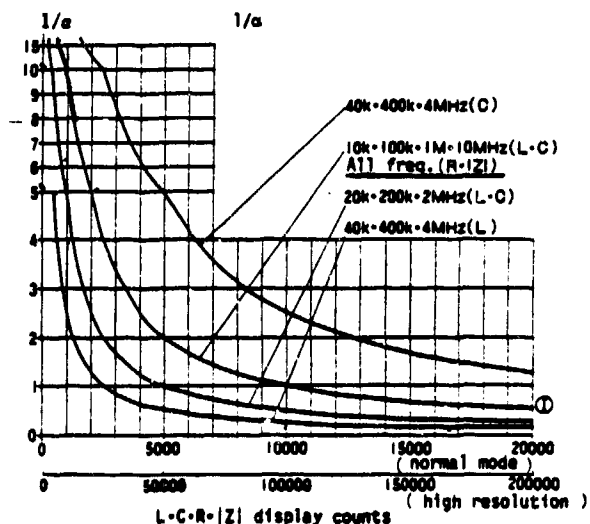
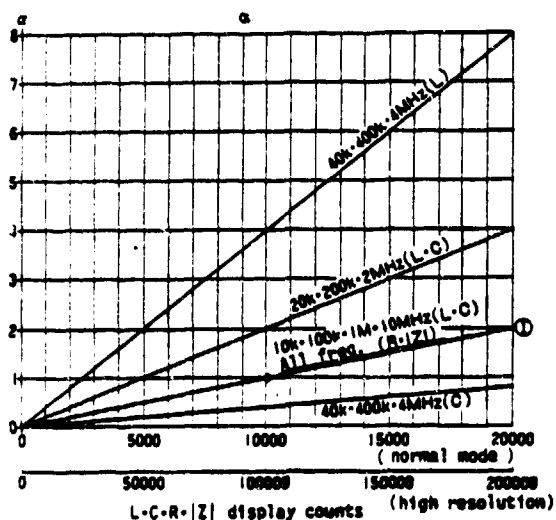


Figure 3 LCR Meter Accuracy

The test cell changes between the open and short circuit measurements and the actual measurement of electrolytic solutions. This is because of the double-layer effect at each electrode and the difference in self-inductance between a good conductor (mercury) and a fair conductor (the electrolytic solutions). These effects cannot be nulled out by the LCR meter since they are not present during the open and short circuit tests. The double layer effect is caused by the electronic-to-ionic conductor interface. The result is the formation of a leaky capacitor at each electrode, as shown in Figure 4. Also, the self-inductance of the test cell is much larger for the electrolytic solutions than for mercury. Calculations for an infinitely long conductor indicate the difference to be around 5 nH. However, end effects will certainly increase this number.

2.2

Results

Measurements on carbon composition resistors were made to test the system background. Since the resistors are electronic conductors, the effects of conductivity relaxation will occur at frequencies much higher than ionic conductors. Two resistors were used, one with a resistance of 228 Ω and the other with a resistance of 7570 Ω . This range of values of resistance is close to those of the electrolytic solutions studied. Tables 1 and 2 list the results of the carbon composition resistors. The format is typical of all experimental results. Background measurements are made after the open- and short-circuit calibration measurements [a zero offset adjustment (ZOA)], separately. Thus, the stray capacitance and residual inductance of the network are recorded. The stray or parasitic background capacitance is measured at 40 kHz and the residual or parasitic background inductance is measured at 10 MHz. The measured background capacitance and inductance as listed are typically the meter's limit of sensitivity. The device to be measured is then connected to the test network and approximate values for the inductance and capacitance are

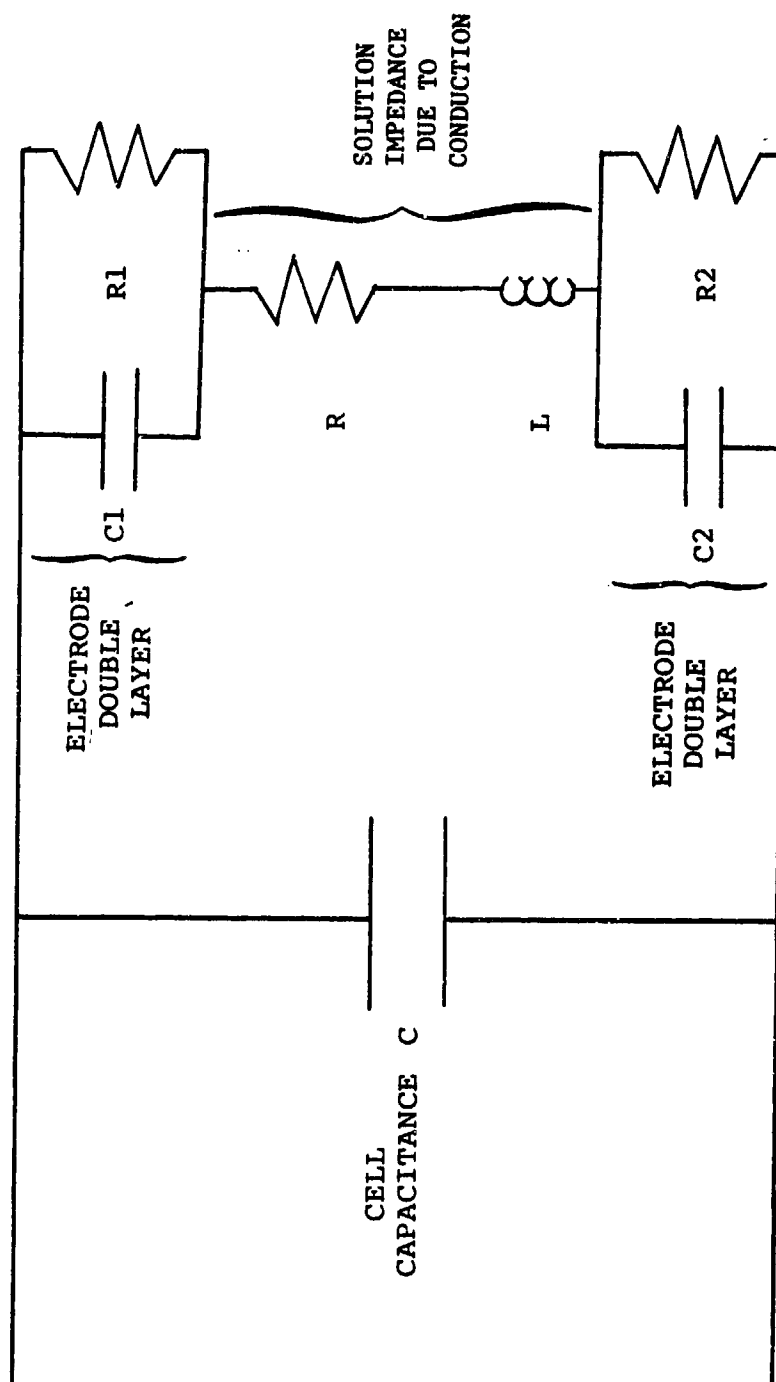


Figure 4 Equivalent Electrical Circuit

Table 1

228Ω CARBON COMPOSITION RESISTOR RESULTS

02/09/84 16:30 CARBON COMP RESISTOR
KEY SETTINGS ARE A1B3C2D0F20H110M2R3180T1

FREQUENCY PARASITIC BACKGROUND CAPACITANCE AND CONDUCTANCE
40K -2.16000000000E-16 .000000000012

FREQUENCY PARASITIC BACKGROUND INDUCTANCE
1GM 5.02600000000E-10

RESISTOR

APPROXIMATE INDUCTANCE
-.000000011086

CELL CAPACITANCE
2.00000000000E-14

FREQUENCY	VOLTAGE	CURRENT	MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
			MAGNITUDE	PHASE	MAGNITUDE	PHASE
1.0E+04	2.40E-01	1.00E-03	2.2828E+02	+0.00	2.2828E+02	+0.00
2.0E+04	2.30E-01	1.00E-03	2.2830E+02	+0.01	2.2830E+02	+0.01
4.0E+04	2.30E-01	1.00E-03	2.2829E+02	+0.00	2.2829E+02	+0.00
1.0E+05	2.30E-01	1.00E-03	2.2830E+02	-0.00	2.2830E+02	-0.00
2.0E+05	2.30E-01	1.00E-03	2.2829E+02	+0.00	2.2829E+02	+0.00
4.0E+05	2.30E-01	1.00E-03	2.2825E+02	+0.01	2.2825E+02	+0.01
1.0E+06	2.30E-01	1.00E-03	2.2825E+02	-0.01	2.2825E+02	-0.01
2.0E+06	2.30E-01	1.00E-03	2.2820E+02	+0.02	2.2820E+02	+0.03
4.0E+06	2.30E-01	1.00E-03	2.2812E+02	+0.03	2.2812E+02	+0.03
1.0E+07	2.10E-01	9.00E-04	2.3046E+02	-0.18	2.3046E+02	-0.16

CHANGED VOLTAGE 16:35

APPROXIMATE INDUCTANCE
-.000000015718

CELL CAPACITANCE
1.40000000000E-13

FREQUENCY	VOLTAGE	CURRENT	MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
			MAGNITUDE	PHASE	MAGNITUDE	PHASE
1.0E+04	2.30E-02	1.00E-04	2.2819E+02	+0.01	2.2819E+02	+0.01
2.0E+04	2.30E-02	1.00E-04	2.2823E+02	+0.01	2.2822E+02	+0.01
4.0E+04	2.30E-02	1.00E-04	2.2823E+02	+0.01	2.2823E+02	+0.01
1.0E+05	2.30E-02	1.00E-04	2.2024E+02	+0.00	2.2824E+02	+0.00
2.0E+05	2.30E-02	1.00E-04	2.2823E+02	+0.01	2.2823E+02	+0.01
4.0E+05	2.30E-02	1.00E-04	2.2819E+02	+0.02	2.2819E+02	+0.02
1.0E+06	2.30E-02	1.00E-04	2.2821E+02	-0.00	2.2821E+02	+0.01
2.0E+06	2.30E-02	1.00E-04	2.2815E+02	+0.01	2.2815E+02	+0.03
4.0E+06	2.20E-02	9.00E-05	2.2809E+02	+0.00	2.2809E+02	+0.05
1.0E+07	2.00E-02	8.00E-05	2.3056E+02	-0.26	2.3056E+02	-0.15

Table 2
75700 CARBON COMPOSITION RESISTOR RESULTS

```

02/09/84 16:04 CARBON COMP RESISTOR
LEV SETTINGS ARE      A1E3C20GF20H110M3R31S0T1
++++
FREQUENCY  PARASITIC BACKGROUND CAPACITANCE AND CONDUCTANCE
40K        -8.600000000000E-17  .0000000000134
++++
FREQUENCY  PARASITIC BACKGROUND INDUCTANCE
160        -1.015800000000E-09
++++

02/09/84 16:04 CARBON COMP RESISTOR
++++
APPROXIMATE INDUCTANCE
-.00000137668
++++

CELL CAPACITANCE
3.636000000000E-13
++++

```

FREQUENCY	VOLTAGE	CURRENT	MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
			MAGNITUDE	PHASE	MAGNITUDE	PHASE
1.0E+04	1.10E-01	0.00E+00	7.5687E+03	+0.07	7.5687E+03	+0.04
2.0E+04	1.10E-01	0.00E+00	7.5707E+03	+0.03	7.5707E+03	+0.05
4.0E+04	1.10E-01	0.00E+00	7.5714E+03	-0.00	7.5714E+03	+0.04
1.0E+05	1.10E-01	0.00E+00	7.5729E+03	-0.11	7.5729E+03	-0.31
2.0E+05	1.10E-01	0.00E+00	7.5701E+03	-0.14	7.5701E+03	+0.06
4.0E+05	1.10E-01	0.00E+00	7.5674E+03	-0.30	7.5675E+03	+0.10
1.0E+06	1.10E-01	0.00E+00	7.5619E+03	-0.84	7.5627E+03	+0.17
2.0E+06	1.10E-01	0.00E+00	7.5395E+03	-1.66	7.5425E+03	+0.35
4.0E+06	1.10E-01	0.00E+00	7.5048E+03	-3.19	7.5157E+03	+0.81
1.0E+07	1.00E-01	0.00E+00	7.4039E+03	-7.11	7.4525E+03	+2.79

obtained. Then the computer directs the LCR meter to take impedance measurements at the ten available frequencies between 10 kHz and 10 MHz. In addition to the frequency and the magnitude and phase of the impedance, the computer lists the approximate applied voltage and current, and a corrected magnitude and phase impedance. The measured cell capacitance is removed from the measured impedance to give the corrected impedance. The cell capacitance in the case of the carbon composition resistors is the capacitance between the two metal contacts at opposite ends of the resistor. In the case of the modified Jones cell, it is typically the capacitance between the two spherical ends of the capillary tube. For the carbon composition resistors, the cell capacitance is very small (0.4 to 0.1 pF or less). The impedance of the carbon composition resistors is largely resistive with a small capacitive influence at the higher frequencies. Notice that the listed approximate inductance is negative. This means the impedance was actually capacitive and not inductive. Thus the LCR meter could not properly interpret the result.

Tables 3, 4, and 5 list the main results of the experiment. For sodium chloride and seawater solutions (concentrations are listed in parts per thousand of solute over solution by weight), the impedance is purely resistive within the accuracy of the LCR meter. The inductive effect at the higher frequencies is less than 0.1 percent. In Table 3, the cell capacitance dominates at all frequencies and only the corrected impedance shows an inductive character because of the high resistance. This result is different from previous less-precise observations that showed a 1- to 2-percent inductive effect at the high frequency end (Reference 4). It is interesting to note that the electrodes in the present experiment have approximately the same surface area and separation as the electrodes in the previous experiment. Because of the capillary tube in the modified Jones cell, however, the cell constant is approximately an order of magnitude larger. If the observed inductance in the previous experiments was caused by some electrode process, it would explain the order of magnitude smaller inductive effect in the present

Table 3

NaCl SOLUTION IMPEDANCE

03/13/84 14:57 NaCl S=13.4PPT T=23.5C
KEY SETTINGS ARE AIB3C2D0F23H110M3P31S0T1

FREQUENCY PARASITIC BACKGROUND CAPACITANCE AND CONDUCTANCE
40K -3.20000000000E-17 -.000000000000

FREQUENCY PARASITIC BACKGROUND INDUCTANCE
10M 6.39500000000E-10

15:01 T=23.4

APPROXIMATE INDUCTANCE
-.0000000331

CELL CAPACITANCE
6.50000000000E-13

			MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
FREQUENCY (HZ)	VOLTAGE (VOLTS)	CURRENT (AMPS)	MAGNITUDE (OHMS)	PHASE (DEG)	MAGNITUDE (OHMS)	PHASE (DEG)
1.0E+04	9.90E-01	1.10E-03	7.7227E+02	-0.04	7.7227E+02	-0.04
2.0E+04	8.90E-01	1.10E-03	7.7229E+02	-0.01	7.7229E+02	-0.01
4.0E+04	8.90E-01	1.10E-03	7.7215E+02	-0.01	7.7215E+02	+0.03
1.0E+05	8.90E-01	1.10E-03	7.7196E+02	-0.01	7.7196E+02	+0.01
2.0E+05	8.90E-01	1.10E-03	7.7169E+02	-0.00	7.7169E+02	+0.03
4.0E+05	8.90E-01	1.10E-03	7.7131E+02	+0.01	7.7131E+02	+0.08
1.0E+06	8.90E-01	1.10E-03	7.7112E+02	-0.04	7.7112E+02	+0.14
2.0E+06	8.80E-01	1.10E-03	7.7063E+02	-0.03	7.7062E+02	+0.33
4.0E+06	8.70E-01	1.10E-03	7.6952E+02	-0.09	7.6947E+02	+0.53
1.0E+07	7.50E-01	9.00E-04	7.7026E+02	-0.15	7.6994E+02	+1.65

15:03 T=23.4

APPROXIMATE INDUCTANCE
-.000000032

CELL CAPACITANCE
5.94000000000E-13

			MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
FREQUENCY (HZ)	VOLTAGE (VOLTS)	CURRENT (AMPS)	MAGNITUDE (OHMS)	PHASE (DEG)	MAGNITUDE (OHMS)	PHASE (DEG)
1.0E+04	9.00E-01	1.10E-03	7.7031E+02	-0.04	7.7031E+02	-0.04
2.0E+04	8.90E-01	1.10E-03	7.7031E+02	-0.01	7.7031E+02	-0.01
4.0E+04	8.90E-01	1.10E-03	7.7022E+02	-0.01	7.7022E+02	+0.00
1.0E+05	8.90E-01	1.10E-03	7.7014E+02	-0.01	7.7014E+02	+0.00
2.0E+05	8.90E-01	1.10E-03	7.7002E+02	-0.01	7.7002E+02	+0.03
4.0E+05	8.90E-01	1.10E-03	7.6900E+02	+0.00	7.6900E+02	+0.07
1.0E+06	8.80E-01	1.10E-03	7.6970E+02	-0.04	7.6970E+02	+0.13
2.0E+06	8.80E-01	1.10E-03	7.6924E+02	-0.03	7.6923E+02	+0.30
4.0E+06	8.70E-01	1.10E-03	7.6821E+02	-0.09	7.6817E+02	+0.57
1.0E+07	7.50E-01	9.00E-04	7.6898E+02	-0.15	7.6872E+02	+1.50

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Table 4

NaCl SOLUTION IMPEDANCE

08/13/84 15:21 NaCl S=50PPT T=23.70

KEY SETTINGS ARE R4B3C2D0F11H11CM3R31S0T1

FREQUENCY PARASITIC BACKGROUND CAPACITANCE AND CONDUCTANCE

40N -2.80000000000E-17 .00000000000

FREQUENCY PARASITIC BACKGROUND INDUCTANCE

10N .000000000052

15:20 T=23.80

APPROXIMATE INDUCTANCE

.000000014152

CELL CAPACITANCE

2.02000000000E-12

				MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
FREQUENCY	VOLTAGE	CURRENT	MAGNITUDE	PHASE	MAGNITUDE	PHASE	
(HZ)	(VOLTS)	(AMPS)	(OHMS)	(DEG)	(OHMS)	(DEG)	
1.0E+04	7.00E-01	3.00E-03	2.3014E+02	-0.11	2.3014E+02	-0.11	
2.0E+04	7.00E-01	3.00E-03	2.3002E+02	-0.05	2.3002E+02	-0.05	
4.0E+04	7.00E-01	3.00E-03	2.2995E+02	-0.03	2.2995E+02	-0.02	
1.0E+05	7.00E-01	3.00E-03	2.2995E+02	-0.02	2.2995E+02	+0.00	
2.0E+05	7.00E-01	3.00E-03	2.2992E+02	-0.01	2.2992E+02	+0.03	
4.0E+05	6.90E-01	3.00E-03	2.2985E+02	+0.01	2.2984E+02	+0.08	
1.0E+06	7.00E-01	3.00E-03	2.2971E+02	+0.00	2.2971E+02	+0.17	
2.0E+06	7.00E-01	3.00E-03	2.2953E+02	+0.05	2.2953E+02	+0.39	
4.0E+06	6.90E-01	3.00E-03	2.2918E+02	+0.07	2.2916E+02	+0.74	
1.0E+07	6.10E-01	2.60E-03	2.2930E+02	+0.22	2.2918E+02	+1.09	

15:31 T=23.70

APPROXIMATE INDUCTANCE

.00000001411

CELL CAPACITANCE

1.92000000000E-12

				MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
FREQUENCY	VOLTAGE	CURRENT	MAGNITUDE	PHASE	MAGNITUDE	PHASE	
(HZ)	(VOLTS)	(AMPS)	(OHMS)	(DEG)	(OHMS)	(DEG)	
1.0E+04	7.00E-01	3.00E-03	2.2957E+02	-0.11	2.2957E+02	-0.10	
2.0E+04	7.00E-01	3.00E-03	2.2951E+02	-0.05	2.2951E+02	-0.05	
4.0E+04	7.00E-01	3.00E-03	2.2947E+02	-0.03	2.2947E+02	-0.02	
1.0E+05	7.00E-01	3.00E-03	2.2945E+02	-0.02	2.2945E+02	+0.00	
2.0E+05	7.00E-01	3.00E-03	2.2942E+02	-0.01	2.2942E+02	+0.00	
4.0E+05	6.90E-01	3.00E-03	2.2937E+02	+0.01	2.2937E+02	+0.07	
1.0E+06	6.90E-01	3.00E-03	2.2933E+02	+0.00	2.2932E+02	+0.16	
2.0E+06	7.00E-01	3.00E-03	2.2923E+02	+0.05	2.2923E+02	+0.37	
4.0E+06	6.90E-01	3.00E-03	2.2894E+02	+0.07	2.2892E+02	+0.71	
1.0E+07	6.10E-01	2.60E-03	2.2909E+02	+0.23	2.2898E+02	+1.01	

Table 5

SIMULATED SEAWATER SOLUTION IMPEDANCE

```

02 16/84 12:59 INSTANT SEA 35 PPT T=23.4
NEW SETTINGS ARE      A4E9C2D0F11H110M3R31S0T1
*****
FREQUENCY  PARASITIC BACKGROUND CAPACITANCE AND CONDUCTANCE
400        6.200000000000E-17  -.0000000000028
*****
FREQUENCY  PARASITIC BACKGROUND INDUCTANCE
10M        6.326000000000E-10

T=23.50
*****
APPROXIMATE INDUCTANCE
.00000001294

CELL CAPACITANCE
1.240000000000E-12
*****

```

FREQUENCY (HZ)	VOLTAGE (VOLTS)	CURRENT (AMPS)	MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
			MAGNITUDE (OHMS)	PHASE (DEG)	MAGNITUDE (OHMS)	PHASE (DEG)
1.0E+04	7.80E-01	2.20E-03	3.4419E+02	-0.05	3.4419E+02	-0.09
2.0E+04	7.80E-01	2.20E-03	3.4417E+02	-0.05	3.4417E+02	-0.04
4.0E+04	7.80E-01	2.20E-03	3.4415E+02	-0.03	3.4415E+02	-0.02
1.0E+05	7.80E-01	2.20E-03	3.4415E+02	-0.02	3.4415E+02	-0.00
2.0E+05	7.80E-01	2.20E-03	3.4413E+02	-0.01	3.4413E+02	+0.02
4.0E+05	7.70E-01	2.20E-03	3.4408E+02	+0.01	3.4407E+02	+0.07
1.0E+06	7.70E-01	2.20E-03	3.4405E+02	-0.01	3.4404E+02	+0.15
2.0E+06	7.70E-01	2.20E-03	3.4392E+02	+0.04	3.4391E+02	+0.34
4.0E+06	7.60E-01	2.20E-03	3.4351E+02	+0.04	3.4349E+02	+0.65
1.0E+07	6.70E-01	1.90E-03	3.4398E+02	+0.13	3.4383E+02	+1.67

```

23.7 C
*****
APPROXIMATE INDUCTANCE
.00000001254

CELL CAPACITANCE
1.000000000000E-12
*****

```

FREQUENCY (HZ)	VOLTAGE (VOLTS)	CURRENT (AMPS)	MEASURED IMPEDANCE		CORRECTED IMPEDANCE	
			MAGNITUDE (OHMS)	PHASE (DEG)	MAGNITUDE (OHMS)	PHASE (DEG)
1.0E+04	7.80E-01	2.20E-03	3.4418E+02	-0.09	3.4418E+02	-0.09
2.0E+04	7.80E-01	2.20E-03	3.4411E+02	-0.04	3.4411E+02	-0.04
4.0E+04	7.80E-01	2.20E-03	3.4405E+02	-0.03	3.4405E+02	-0.02
1.0E+05	7.80E-01	2.20E-03	3.4402E+02	-0.02	3.4402E+02	-0.00
2.0E+05	7.60E-01	2.20E-03	3.4397E+02	-0.01	3.4397E+02	+0.03
4.0E+05	7.70E-01	2.20E-03	3.4389E+02	+0.01	3.4388E+02	+0.07
1.0E+06	7.70E-01	2.20E-03	3.4384E+02	-0.01	3.4384E+02	+0.15
2.0E+06	7.70E-01	2.20E-03	3.4369E+02	+0.04	3.4368E+02	+0.36
4.0E+06	7.60E-01	2.20E-03	3.4327E+02	+0.04	3.4324E+02	+0.68
1.0E+07	6.70E-01	1.90E-03	3.4372E+02	+0.13	3.4356E+02	+1.74

experiment. That is, the electrode process remained unchanged while the cell impedance was increased by an order of magnitude. Further, the relative changes in the impedance are essentially the same in the previous and present experiments.

Another set of measurements attempted to observe the Debye-Falkenhagen effect (Reference 3) that is observable within the frequency band of the LCR meter but at very low concentrations (around 1000 times less than seawater concentrations). Tables 6 and 7 show the results for 0.01-Molarity (M) solutions. The Debye-Falkenhagen effect is much stronger for MgSO_4 than for KCl ; thus, a comparison can be made between the two tables. Within the accuracy of the measurements, no differences are observed and the corrected impedance is purely resistive. The concentration is too high to observe this effect.

Finally, much of the initial motivation to carry out these experiments was to measure the conductivity relaxation time. This number is at least 10^{-13} sec, which is based on the initial assumptions of Debye and Falkenhagen (Reference 3). Based on the experimental results of Table 5, the maximum seawater conductivity relaxation time, τ_c , is 5×10^{-10} sec. Therefore,

$$1 \times 10^{-13} < \tau_c < 5 \times 10^{-10} \text{ sec or}$$

$$3.2 \times 10^8 < f_c \left(= \frac{1}{2\pi\tau_c} \right) < 1.6 \times 10^{12} \text{ Hz}$$

3. PRESENT UNDERSTANDING OF THE ELECTRICAL PROPERTIES OF SEAWATER

The problem of understanding the electrical properties of seawater is a very complex one. This is because the detailed physics of electrolytic solutions is not completely understood and seawater is a heterogeneous solu-

Table 6
KC1 SOLUTION IMPEDANCE

02/10/84 14:43 KCL .012M T=23.0									
AL7 SETTINGS ARE RAJ30260F11H110N3P01S011									
+++++									
FREQUENCY PARASITIC BACKGROUND CAPACITANCE AND CONDUCTANCE									
400									
-6.00000000000E-17 .00000000000E+00									
+++++									
FREQUENCY PARASITIC BACKGROUND INDUCTANCE									
100									
6.50000000000E-10									
+++++									
14:54 T=23.8									
APPROXIMATE INDUCTANCE									
-0.00000000000E+00									
+++++									
CELL CAPACITANCE									
2.19200000000E-13									
+++++									
FREQUENCY (HZ)	VOLTAGE (VOLTS)	CURRENT (AMPS)	MEASURED MAGNITUDE (OHMS)	IMPEDANCE PHASE (DEG)	CORRECTED MAGNITUDE (OHMS)	IMPEDANCE PHASE (DEG)			
1.0E+04	1.00E+00	0.00E+00	1.2899E+04	+0.01	1.2899E+04	+0.02			
2.0E+04	1.00E+00	0.00E+00	1.2901E+04	+0.00	1.2901E+04	+0.02			
4.0E+04	1.00E+00	0.00E+00	1.2900E+04	-0.01	1.2900E+04	+0.03			
1.0E+05	1.00E+00	0.00E+00	1.2899E+04	-0.11	1.2899E+04	-0.01			
2.0E+05	1.00E+00	0.00E+00	1.2895E+04	-0.17	1.2895E+04	+0.04			
4.0E+05	1.00E+00	0.00E+00	1.2889E+04	-0.33	1.2889E+04	+0.08			
1.0E+06	9.90E-01	0.00E+00	1.2885E+04	-0.86	1.2886E+04	+0.16			
2.0E+06	9.90E-01	0.00E+00	1.2878E+04	-1.65	1.2883E+04	+0.38			
4.0E+06	9.80E-01	0.00E+00	1.2899E+04	-3.34	1.2920E+04	+0.73			
1.0E+07	8.20E-01	0.00E+00	1.3893E+04	-9.12	1.3257E+04	+1.27			

Table 7
MgSO4 SOLUTION IMPEDANCE

```

02/10/84 14:24 MgSO4 .01N I= 23.4C
KEY SETTINGS ARE      A4B3C2D0F11H110M3R31S0T1
+++++
FREQUENCY PARASITIC BACKGROUND CAPACITANCE AND CONDUCTANCE
40K      -6.8000000000E-17 -.000000000002
+++++
FREQUENCY PARASITIC BACKGROUND INDUCTANCE
10K      .000000000000973
+++++

14:32 I=23.6
+++++
APPROXIMATE INDUCTANCE
-.0000238102
+++++

CELL CAPACITANCE
2.03400000000E-13
+++++

```

MEASURED IMPEDANCE		CORRECTED IMPEDANCE			
FREQUENCY (HZ)	VOLTAGE (VOLTS)	CURRENT (AMPS)	PHASE (DEG)	MAGNITUDE (OHMS)	PHASE (DEG)
1.0E+04	1.01E+00	0.00E+00	+0.01	1.1209E+04	+0.02
2.0E+04	1.00E+00	0.00E+00	+0.01	1.1230E+04	+0.02
4.0E+04	1.00E+00	0.00E+00	-0.01	1.1249E+04	+0.02
1.0E+05	1.00E+00	0.00E+00	-0.09	1.1267E+04	-0.01
2.0E+05	1.00E+00	0.00E+00	-0.13	1.1282E+04	+0.04
4.0E+05	1.00E+00	0.00E+00	-0.27	1.1297E+04	+0.07
1.0E+06	9.90E-01	0.00E+00	-0.72	1.1314E+04	+0.13
2.0E+06	9.90E-01	0.00E+00	-1.42	1.1331E+04	+0.23
4.0E+06	9.80E-01	0.00E+00	-2.91	1.1373E+04	+0.50
1.0E+07	8.30E-01	0.00E+00	-7.85	1.1634E+04	+0.83

tion of many electrolytes. The importance of the electrical properties (i.e., dielectric and conductive) at radar frequencies has required some form of tabulation as a function of temperature, salinity, and frequency. This is currently done in a semi-empirical way using the formulas of Debye, and Cole and Cole (References 5 and 6) fitted to experimental data (Reference 7). The result is a reasonable representation of the true dielectric and conductive values of seawater as a function of temperature and salinity below 5 GHz. However, precision microwave radiometry requires equal precision of the known electrical properties of seawater. To improve the present semi-empirical formulas requires a greater understanding of the theory, including observed anomalous electrical properties. A brief review of the present understanding of the electrical properties of seawater and seawater-related electrolytes is presented in the following section.

3.1 Microwave Electrical Properties

Microwave data of the complex relative permittivity for electrolytic solutions is typically fitted to an equation of the form (Reference 7)

$$\epsilon(\omega, T, S) = \epsilon_{\infty} + \frac{\epsilon_s(T, S) - \epsilon_{\infty}}{1 + [j\omega\tau_D(T, S)]^{1-\alpha(T, S)}} - j \frac{\sigma(T, S)}{\omega\epsilon_0} \quad (8)$$

where

T = temperature

S = salinity

ϵ_{∞} = the relative dielectric constant at $\omega \gg 1/\tau_D$

ϵ_s = the static relative dielectric constant

ϵ_0 = permittivity of free space [Farads/m]

$\omega = 2\pi f$, f = frequency in hertz

σ = conductivity [S/m]

τ_D = Debye relaxation time

α = empirical spreading parameter of relaxation times

The conductivity relaxation process is not represented by Equation (8) and neither are higher frequency resonances which impact ϵ_{∞} . However, since this model only claims agreement below 5 GHz, the details of the Debye and conductivity relaxations are not required. For example, for $\alpha = 0$ Equation (9) can be expanded to the form

$$\begin{aligned}\epsilon(\omega, T, S) &= \epsilon_{\infty} + (\epsilon_s - \epsilon_{\infty}) (1 - (j\omega\tau_D)^2 - (\omega\tau_D)^2 \dots) - j \frac{\sigma}{\omega\epsilon_0} \quad (9) \\ &= \epsilon_s - (\epsilon_s - \epsilon_{\infty}) (\omega\tau_D)^2 - j[(\epsilon_s - \epsilon_{\infty})(\omega\tau_D) + \frac{\sigma}{\omega\epsilon_0}]\end{aligned}$$

Because $\omega\tau_D \ll 1$ below 5 GHz, this expansion will work very well. The expansion of any proper function would produce the same first-order structure. Thus, curve fitting data to Equation (8) at frequencies below 5 GHz will not definitively determine the true structure of the dielectric properties beyond 5 GHz. This fact is recognized by Klein and Swift (Reference 7).

In principle, ϵ_{∞} and σ are functions of temperature, salinity, and frequency, and ϵ_s and τ_D are functions of temperature and salinity. A more general form of Equation (9) is given by

$$\begin{aligned}\epsilon(\omega, T, S) &= \epsilon_{\infty}(\omega, T, S) + \frac{\epsilon_s(T, S) - \epsilon_{\infty}(\omega, T, S)}{1 + [j\omega\tau_D(T, S)]^{1-\alpha(T, S)}} \quad (10) \\ &\quad - j \frac{1}{\omega\epsilon_0} \frac{\sigma_{DC}(T, S)}{1 + j\omega\tau_C(T, S)}\end{aligned}$$

where

$$\begin{aligned}\sigma_{DC} &= \text{DC conductivity} \\ \tau_c &= \text{conductivity relaxation time}\end{aligned}$$

The Debye relaxation process accounts for the damped rotation of the water molecule in the liquid state. The conductivity relaxation process accounts for the damped translational motion of the ionic charge carrier. Debye in 1913 (Reference 8) approximately determined the Debye relaxation time, τ_D , by considering Stoke's law converted for a rotating sphere and Brownian motion. The following formula was obtained.

$$\tau_D = \frac{8\pi\eta a^3}{kT} \quad (11)$$

where

$$\begin{aligned}\eta &= \text{viscosity of water} \\ a &= \text{radius of water molecule} \\ k &= \text{Boltzmann's constant}\end{aligned}$$

Surprisingly, this formula predicts reasonably well the observed relaxation time of water, which is on the order of 1×10^{-11} sec. Intuitively, the relaxation time of the translational motion of the ionic charge carriers in water should be similar to the rotational relaxation of the water molecule. Following a similar approach as Debye, the conductivity relaxation time can be approximately found by examining the charge density in light of Stoke's law and Brownian motion. The result is

$$\tau_C = \frac{6\pi\eta b x_o^2}{kT} \quad (12)$$

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where

b = radius of charge carrier

x_0 = $1/e$ point of Maxwell-Boltzmann distribution
of charge density

Since

$$\begin{aligned} b &\sim a \\ x_0 &\geq a \end{aligned}$$

then

$$\tau_c \geq \tau_D$$

This result crudely verifies our intuition. It is interesting to note that the Debye-Falkenhagen computation of the conductivity relaxation time is based on the velocity or mobility part of the conductivity and does not include the random motion of the ion.

The real and imaginary parts of the general complex permittivity [Equation (10)] are for $\alpha = 0$, thus

$$\epsilon = \epsilon' - j\epsilon'' \quad (13)$$

$$\epsilon' = \epsilon_\infty + \frac{\epsilon_S - \epsilon_\infty}{1 + (\omega\tau_D)^2} - \frac{\sigma_{DC} \tau_C / \epsilon_0}{1 + (\omega\tau_C)^2} \quad (14)$$

$$\epsilon'' = \frac{(\epsilon_S - \epsilon_\infty)\omega\tau_D}{1 + (\omega\tau_D)^2} - \frac{\sigma_{DC} \tau_C / \epsilon_0 \omega}{1 + (\omega\tau_C)^2} \quad (15)$$

Figures 5a and 5b are plots of ϵ' and ϵ'' , respectively, for different values of τ_c . The cases for $\tau_c = \tau_D$ or 0 are the same curves. Interestingly, with $\tau_c = \tau_D$ the conductivity effects cannot be differentiated from the case of $\tau_c = 0$. However, for $\tau_c/\tau_D = 0.5$ and 2 there are observable differences. If $\tau_c \sim \tau_D$ (it is unlikely that they are exactly equal) then the conductivity relaxation effect is an important term in the complex permittivity formula. The efforts to model the permittivity beyond 5 GHz would then be helped by including this effect.

Figures 6a and 6b illustrate the dielectric properties of pure water. Figure 6a also shows typical skin depth attenuation of seawater (References 9 and 10). The pure water spectra show a fair amount of structure. The hump at 20 GHz is the Debye relaxation phenomenon. The two peaks between 10^{12} and 3×10^{13} Hz are called the hindered translation, ν_T , and librational band, ν_L , respectively. They are caused by H_2O and H_2O (i.e., dimmer-like) interactions (Reference 11). Most of the higher frequency peaks are caused by the internal motion of the water molecule (i.e., rotation vibration (Reference 11)). Beyond the Debye relaxation, the complex permittivity is represented by ϵ_∞ . It is dominated by these infrared resonances. This can be seen by examining the real part of $\epsilon_\infty (= \epsilon'_\infty - j\epsilon''_\infty)$;

$$4.9 > \epsilon'_\infty > 1.69 > 1$$

Microwave Optical

ϵ_∞ is determined by resonances and relaxations beyond the frequency of interest. Based on the above relationship, the microwave region is heavily influenced by the infrared resonances.

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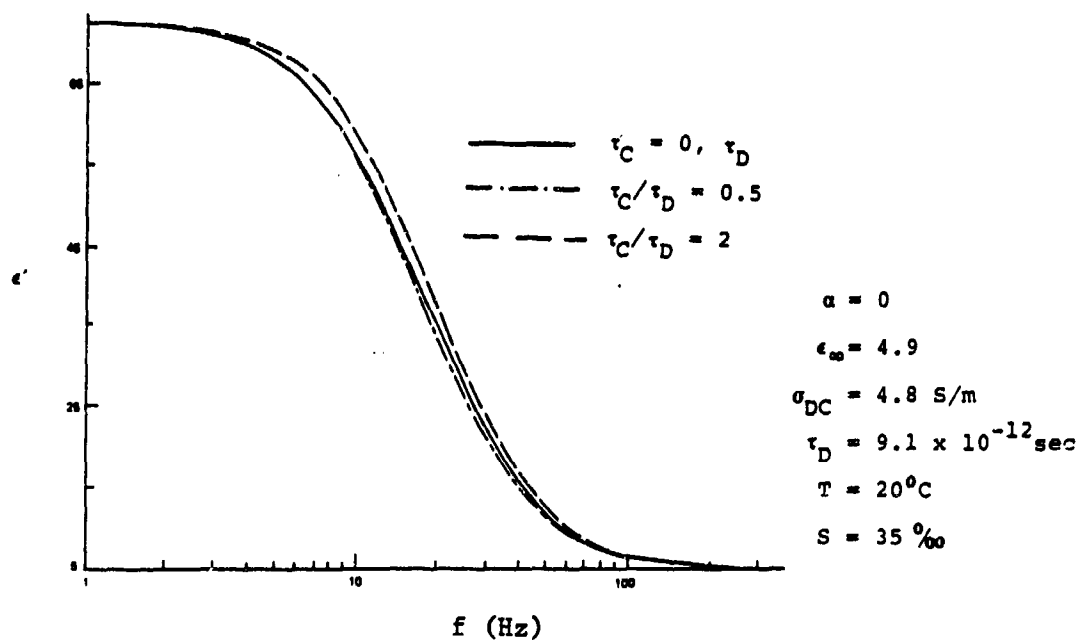


Figure 5(a) Real Part of the Permittivity Versus Frequency

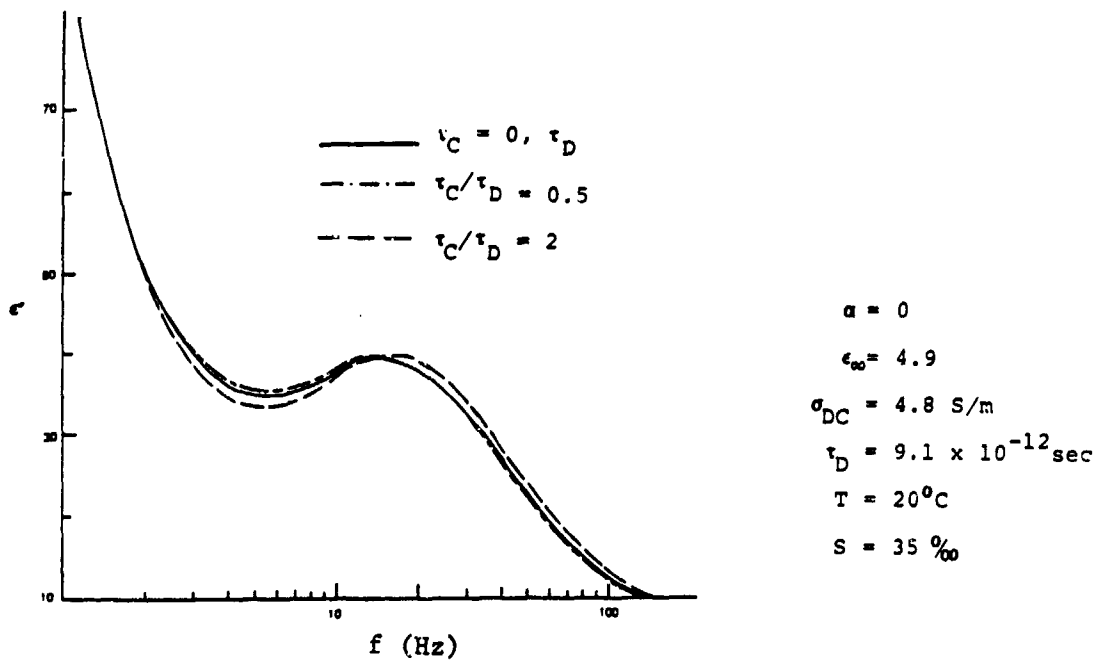
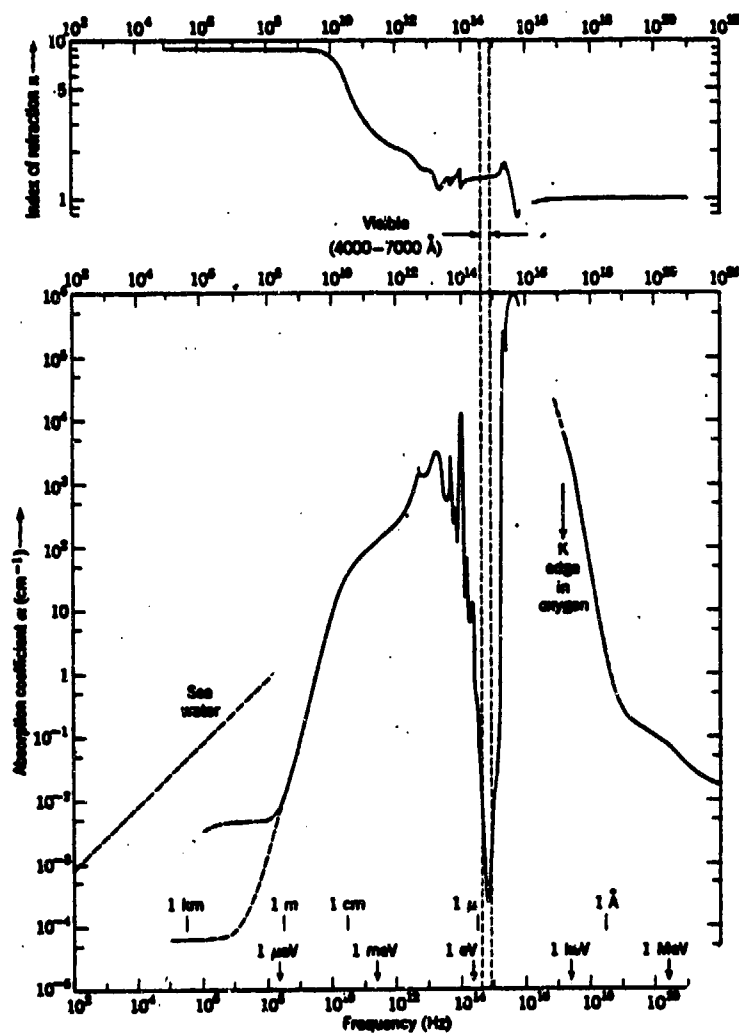


Figure 5(b) Imaginary Part of the Permittivity Versus Frequency

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Note: Also shown as abscissas are an energy scale (arrows) and a wavelength scale (vertical lines). The visible region of the frequency spectrum is indicated by the vertical dashed lines. The absorption of coefficient for seawater is indicated by the dashed diagonal line at the left. Scales are logarithmic in both directions.

Figure 6(a) The Index of Refraction (top) and Absorption Coefficient (bottom) for Liquid Water as a Function of Linear Frequency

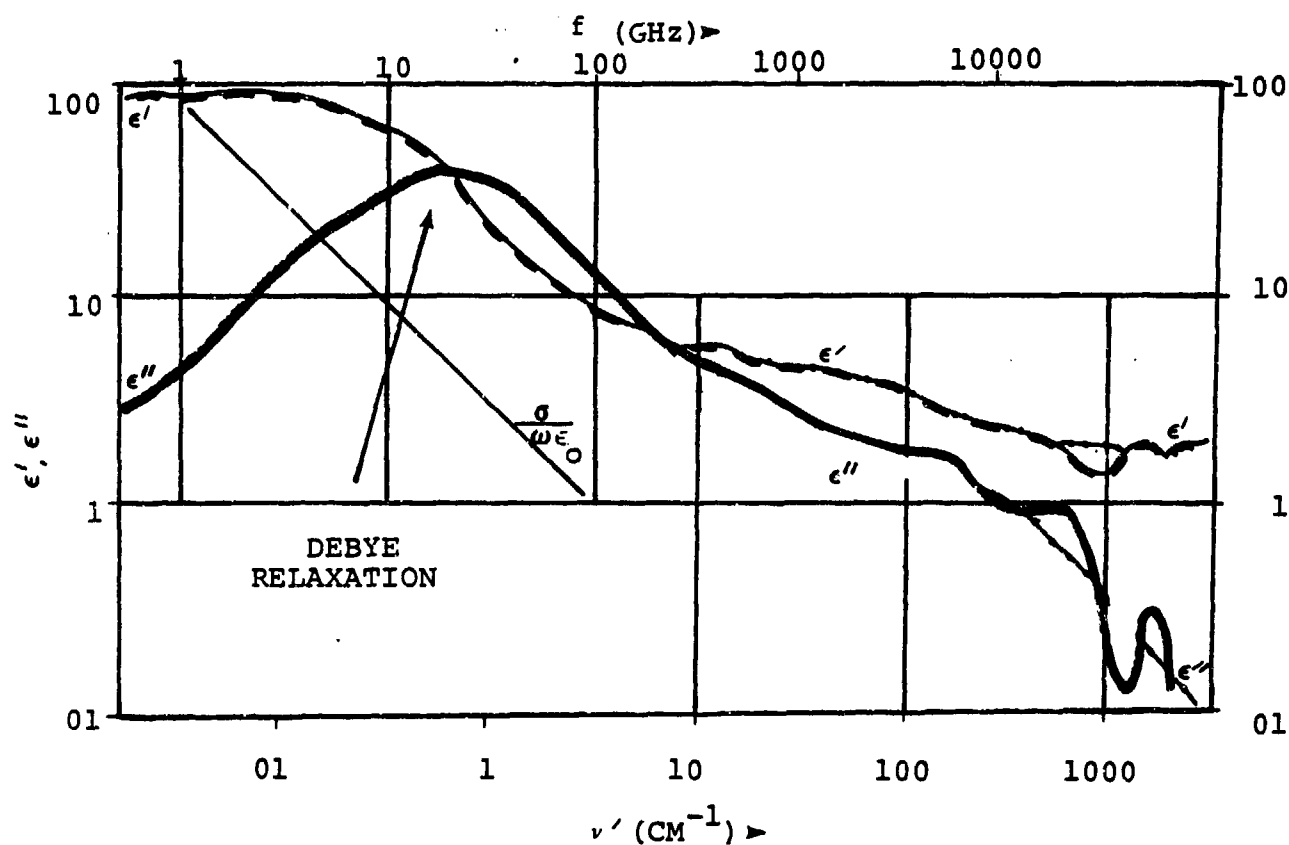


Figure 6(b) Logarithmic Plot of ϵ' (dashed line), ϵ'' (solid line) and $\sigma / \omega \epsilon_0$ (solid line) Against Frequency ν (wave-number ν') for 20°C and $\sigma = 5$ S/m

The presence of a solute does influence ϵ_{∞} . This is demonstrated by Figure 7, which shows experimental data of the ν_L and ν_T band of pure D_2O and D_2O KCl and NaCl solutions at 4 Molarity (M). The solutions are highly concentrated compared to seawater but do show significant changes in the spectrum. This will directly impact ϵ_{∞} at microwave frequencies and make it a function of salinity. It is also well known (Reference 11) that these bands are sensitive to temperature. Thus ϵ_{∞} is indeed a function of frequency, salinity, and temperature.

The relaxation time spreading parameter, α , is not zero and is thought to be a function of temperature and salinity (References 12 and 13). This is illustrated in Figure 8 for H_2O solutions of NaCl. The precise nature of α will not be understood until the conductivity relaxation effects are understood. The relaxation time spreading parameter influences the falloff of the Debye relaxation. As illustrated by Figures 5a and 5b, the conductivity relaxation process also influences the falloff of the curves. Thus, these two different yet similar effects may be difficult to separate.

3.2 Anomalous Electrical Properties

In addition to well-understood physical mechanisms that impact seawater electrical properties, there are reported observations of anomalous behavior of solutions of single electrolytics which compose seawater. These are listed below (Reference 13).

- (1) Anomalously high conductivity of electrolytic solutions at concentrations close to seawater at 582 MHz (Reference 14).
- (2) Additional relaxation phenomena at submillimeter wavelengths. The dielectric constant is modelled as

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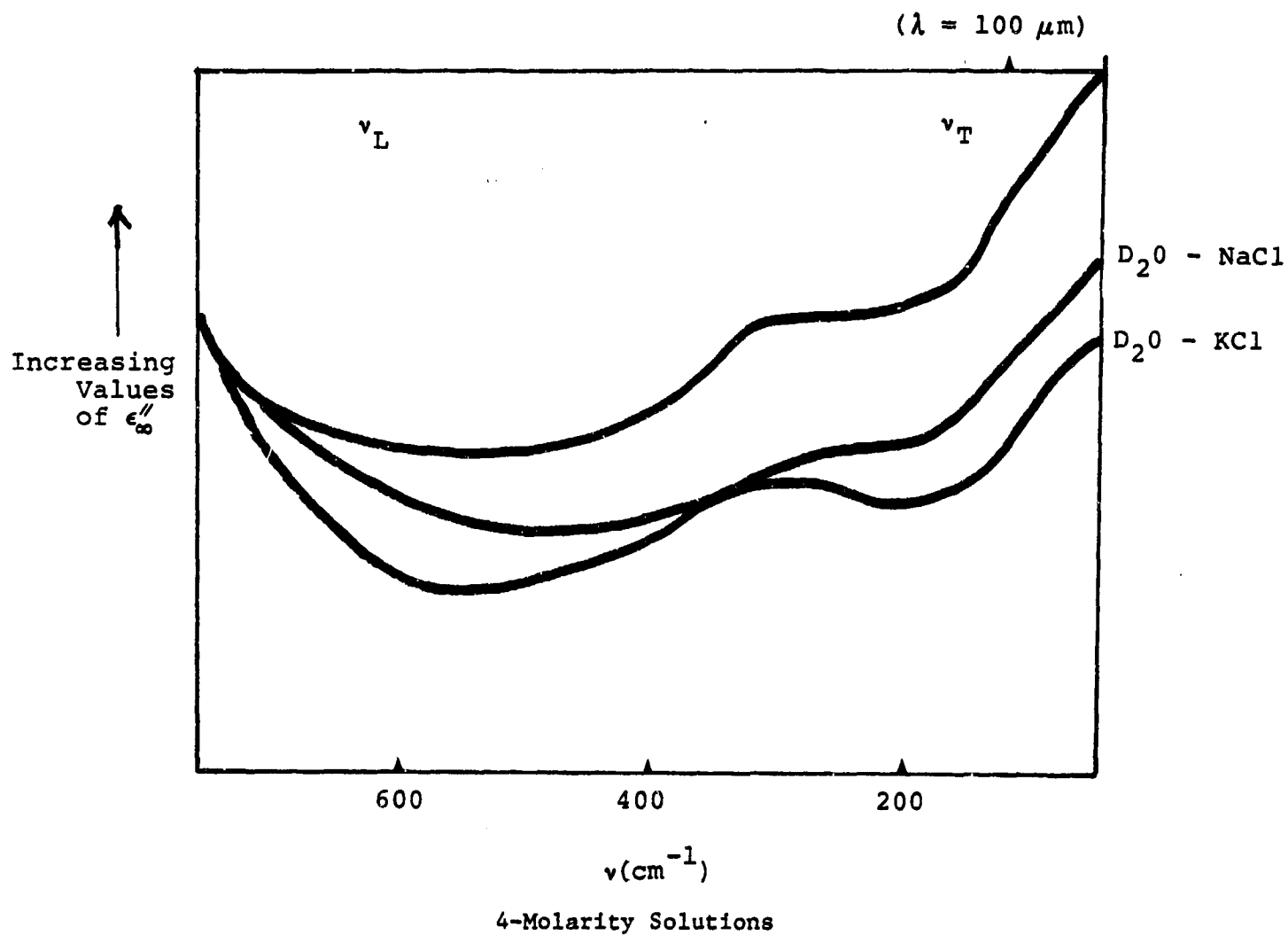


Figure 7 Experimental Data of the ν_L and ν_T Band of Pure D_2O and D_2O KCl and NaCl Solutions at 4-Molarity Concentration

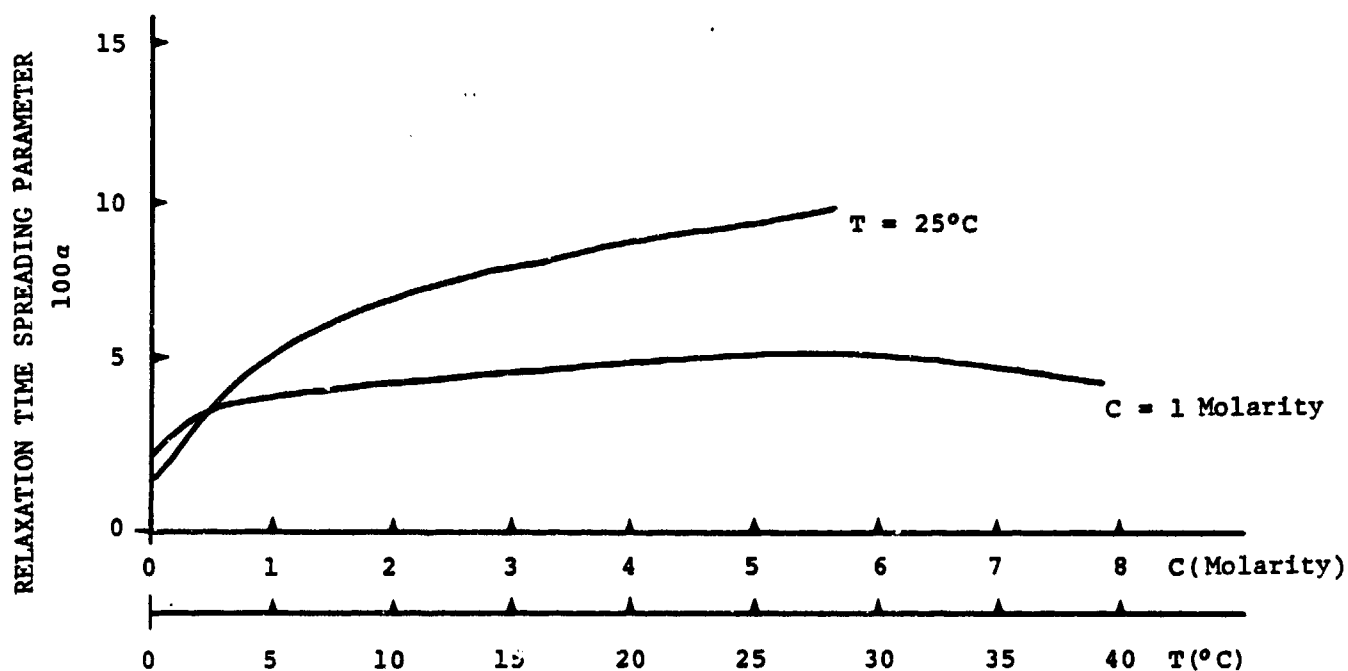


Figure 8 Relaxation Time Spreading Parameter, α , as a Function of Temperature (constant concentration $c = 1 \text{ Molarity}$) and Concentration (constant temperature $T = 25^\circ\text{C}$) for a Water-NaCl Solution

The Johns Hopkins University
APPLIED PHYSICS LABORATORY
Laurel, Maryland

$$\epsilon = n^2 + \frac{\epsilon_{\infty} - n^2}{1 + j\omega\tau_2} \frac{\epsilon_s - \epsilon_{in}}{1 + (j\omega\tau)^{1-\alpha}}$$

where τ_2 is the additional relaxation time.

4. CONCLUSIONS

The electrical properties of seawater have been examined with particular emphasis on the effects of relaxation of the charge carriers with the frequency of the applied field. Such effects have not been definitively quantified either experimentally or theoretically. Experiments conducted were not conclusive. It is possible that the conductivity relaxation time (the characteristic time it takes a translating charge carrier to relax back to equilibrium) is similar to the Debye relaxation time (the characteristic time it takes a rotating molecule to relax back to equilibrium in a fluid or solid), making direct observation of this effect difficult. The conductivity relaxation time has been constrained to be within $5 \times 10^{-10} > \tau_c > 1 \times 10^{-13}$ sec. It is quite possible that $\tau_c \sim \tau_D$ since the dampening mechanisms should be similar. If this is true, the dielectric models will need to include the effects of conductivity relaxation. Also, it is recognized that resonant far-infrared (including submillimeter) bands can influence the dielectric properties at and above 40 GHz.

A more definitive impedance experiment attempting to measure the conductivity relaxation time will require the following:

- A test cell with a much larger cell constant
- Impedance measurements at frequencies close to 1 GHz
- A high degree of precision or a means of nulling the DC resistance of the electrolytic solution

Also, the development of a reasonable theory may aid in the selection of a more optimal temperature and salinity of the samples.

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